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缝洞型油藏井钻遇大尺度部分充填溶洞数学模型

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摘要: 缝洞型油藏不同介质间流体窜流的研究对于大尺度溶洞中原油是否能够得到有效开发具有重大意义。基于缝洞型油藏大尺度溶洞充填特征,建立了井钻遇大尺度部分充填溶洞数学模型,采用 Laplace 变换和 Stehfest 数值反演,分别得到了基质—溶洞未充填区域窜流量、溶洞充填区域—溶洞未充填区域窜流量和大尺度溶洞无因次流量,并分析了不同参数对窜流特征曲线的影响。研究表明:流体窜流过程可划分为4个流动阶段,流动前期和中前期主要为基质中流体向溶洞未充填部分窜流;流动中后期和后期主要为溶洞充填物流体向溶洞未充填部分窜流。在流动前期,基质和未充填溶洞间流体交换对部分充填溶洞流量贡献较大;而流动后期,溶洞充填物和未充填溶洞间流体交换对部分充填溶洞流量贡献较大。重力会导致溶洞充填物—溶洞未充填部分窜流量减小,而基质—溶洞未充填部分窜流不受到重力影响。溶洞充填程度、溶洞未充填部分和基质系统能量等因素对窜流特征曲线具有重大的影响。研究方法和结果对合理分析缝洞型油藏大尺度溶洞流动特征具有一定的指导意义。

关键词: 缝洞型油藏;大尺度溶洞;充填特征;窜流特征曲线;参数分析;石油地质。

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Mathematical Model for Wells Drilled in Large-Scale Partially Filled Cavity in Fractured-Cavity Reservoirs

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Abstract: It is significant to explore fluid crossflow rule of various medium in fractured-cavity reservoirs for effective development of crude oil in large-scale partially filled cavity. Based on the filling characteristics of the large size cavity in fractured-cavity reservoir, a fluid flow mathematical model for wells drilled in large-scale partially filled cavity in fractured-cavity reservoir has been proposed in this paper. The Laplace transformation and Stehfest numerical inversion are applied to analyze the crossflow rule in the partially filled cavity. The fluid crossflow characteristic curve for the partially filled cavity has been obtained and studied in this paper. The results show that the process of the flow for the partially filled cavity can be divided into four sections, including the earlier period and middle-early period flow for the crossflow between the matrix system and unfilled part of the cavity, the middle-late period and the later period flow for the crossflow between the filler of the cavity and unfilled part of the cavity. In middle-early flow period, the crossflow between the matrix system and unfilled part of the cavity reaches the quasi steady state. And the crossflow between the filler of the cavity and unfilled part of the cavity reaches the quasi steady state at the later period. Gravity is one of the main factors that influences the typical crossflow characteristic curve, leading to the decrease of the crossflow between the filler and the unfilled part of the cavity. However, the inter-porosity flow between the matrix and the unfilled part of the cavity has no relation with gravity. In addition, the cross flow characteristic curve is greatly affected by filling degree of the cavity, the energy parameters of the filler of the cavity and the energy parameters of the matrix. The methods and results of this research can facilitate further studies on the crossflow characteristics of the partially filled cavity.

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ty in fractured-cavity reservoirs.

Key words: fractured-cavity reservoir; large-scale cavity; filling characteristics; cross flow characteristic curve; parameter analysis; petroleum geology.

塔河缝洞型油藏储集空间复杂,裂缝、大尺度溶洞并存,溶洞—裂缝具有空间配置关系复杂且充填类型、充填程度多样等特点,该油藏地层属于空间离散介质,油藏开发存在巨大困难(韩剑发等,2007; Popov *et al.*, 2007; 刘中春等,2009; 李阳和范智慧,2011; 李阳,2012, 2013; 潘建国等,2012; 蔡建超等,2013; 钱海涛等,2014; 王敬等,2014; Cheng *et al.*, 2015; 霍志鹏等,2016).前人研究表明由于大尺度溶洞的存在,三重连续介质渗流模型并不适用于研究塔河缝洞型油藏流动规律(康志江,2010; 李鹞和李允,2010; 姚军等,2010, 2013; 王敬等,2012).国内外不少学者将岩块和裂缝系统视为渗流区域,将溶洞系统视为自由流动区域,建立流态耦合模型对缝洞型油藏渗流规律进行了研究(Arbogast and Gomez, 2008; Peng *et al.*, 2009; 姚军等,2010).熊伟等(2011)将溶洞视为扩大井筒,建立单井钻遇孤立溶洞模型,研究了基岩与孤立溶洞间的不稳定窜流规律.张福祥等(2009)和陈方方等(2015)建立了钻井打在大溶洞内的试井解释数学模型,给出了 Laplace 空间解,并对影响井底压力动态的主要因素进行了分析.刘洪等(2012)将溶洞视为等势体,利用溶洞质量守恒原理建立了含大尺度溶洞缝洞型油藏数学模型,使用直接边界元方法计算了溶洞压力响应,得到了不同参数对溶洞压力响应曲线的影响.

大尺度溶洞(洞穴)一般充填着疏松充填物质(如:砂、碎石和泥)(Popov *et al.*, 2007; 韩剑发等,2007; 潘建国等,2012; 胡向阳等,2014; 钱海涛等,2014; 金强等,2015),溶洞的充填程度是影响储层物性和流动规律的重要因素,而充填物中流体向溶洞未充填区域窜流规律对于缝洞型油藏大尺度溶洞中原油是否能够得到有效开发具有重大意义.而基质和未充填溶洞间流体交换也会对缝洞型油藏流体流动规律产生重大影响.目前,国内外学者对于不同缝洞介质间流体窜流规律还未进行深入研究(程倩等,2009; Gao *et al.*, 2016).为此,笔者根据缝洞型油藏大尺度溶洞充填特征,建立了缝洞型油藏井钻遇大尺度部分充填溶洞(洞穴)数学模型,采用 Laplace 变换和 Stehfest 数值反演,得到了大尺度溶洞窜流特征曲线,分析了各类参数对窜流量的敏感性.

1 模型的建立

1.1 物理模型

根据缝洞型油藏大尺度溶洞或洞穴(组构要素类型主要包括地下河、厅堂洞和竖井洞,特征尺度可达米以上)充填特征(李阳和范智慧,2011; 胡向阳等,2014; 金强等,2015),部分充填溶洞可抽象为如图 1a 所示的物理模型,即溶洞由 2 个半径为 r_c 同心圆柱形区域组成.溶洞周围为具有一定渗流能力的基岩(基岩储渗能力主要由微裂缝贡献),如图 1b 所示.将溶洞充填区域及周围基岩等效为均匀介质,则部分充填溶洞流量主要由溶洞未充填区域周围基岩及等效均匀介质两部分贡献.油藏边界为有限大封闭边界,半径为 r_e ,溶洞未充填区域周围基质渗透率为 k_m ,孔隙度为 φ_m ,综合压缩系数为 C_{tm} ,等效均匀介质渗透率为 k_1 ,孔隙度为 φ_1 ,综合压缩系数为 C_{t1} ,油藏有效厚度为 h ,溶洞未填充部分厚度为 h_1 ,溶洞填充部分厚度为 h_2 .

1.2 数学模型

根据物理模型,以溶洞充填区域—未充填区域交界面处为坐标原点,建立如图 1a 空间坐标系,考虑基岩发生平面径向流动,建立如下数学模型.

$$\begin{cases} \frac{\partial^2 p_m}{\partial r^2} + \frac{1}{r} \frac{\partial p_m}{\partial r} = \frac{\varphi_m t C_{tm}}{k_m} \frac{\partial p_m}{\partial t}; \\ \frac{\partial^2 p_1}{\partial z^2} = \frac{\varphi_1 t C_{t1}}{k_1} \frac{\partial p_1}{\partial t}; \\ p_m|_{t=0} = p_1|_{t=0} = p_v|_{t=0} = p_i; \left. \frac{\partial p_m}{\partial r} \right|_{r=r_e} = 0; \\ p_m|_{r=r_c} = p_v|_{r=r_c}; p_1|_{z=0} = p_v|_{z=0}; \left. \frac{\partial p_1}{\partial z} \right|_{z=h_1} = 0; \\ \frac{2\pi r h_1 k_m}{\mu} \left. \frac{\partial p_m}{\partial r} \right|_{r=r_c} - \frac{\pi r_c^2 h_1 C_L}{\mu} \left. \frac{\partial p_v}{\partial t} \right|_{r=r_c} + \\ \frac{\pi r_c^2 k_1}{\mu} \left(\frac{\partial p_1}{\partial z} - \rho g \right) \Big|_{z=0} - \pi r_c^2 h_1 C_L \left. \frac{\partial p_v}{\partial t} \right|_{z=0} = qB; \end{cases} \quad (1)$$

其中, p 为地层压力, MPa; p_i 为原始地层压力, MPa; C_{tm} 为基质综合压缩系数, MPa^{-1} ; r 为径向距离, m; r_c 为溶洞半径, m; k 为渗透率, μm^2 ; h 为地层有效厚度, m; h_1 为溶洞未充填区域厚度, m; h_2 为溶洞充填区域厚度, m; r_e 为地层半径, m; C_{t1} 为等效均匀介质综合压缩系数, MPa^{-1} ; C_L 为溶洞未充填区域流体压缩系数, MPa^{-1} ; t 为生产时间, d; φ

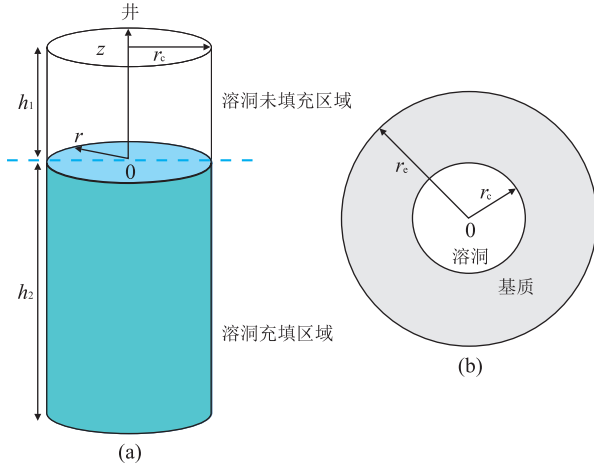


图 1 溶洞-基质型缝洞油藏示意

Fig.1 The sketch of large size vug and matrix in fractured-vuggy reservoir

为孔隙度,%; μ 为流体黏度, $\text{mPa} \cdot \text{s}$; q 为油井产量, m^3/d ; B 为流体体积系数; ρ 为流体密度, kg/m^3 ; g 为重力加速度, N/kg ; 下标: m 为基质; v 为溶洞未充填区域; 1 为等效均匀介质. 公式(1)为缝洞型油藏井钻遇大尺度部分充填溶洞定产量生产时控制方程及初始边界条件. 其中公式(1)的最后一部分包括 4 项: 左边第 1 项为基质-溶洞未充填区域窜流量; 左边第 2 项为溶洞未充填区域对于基质系统弹性释放量; 左边第 3 项为等效均匀介质-溶洞未充填区域窜流量; 左边第 4 项为溶洞未充填区域对于等效均匀介质弹性释放量.

1.3 模型求解

引入如下无量纲量:

$$p_D = \frac{2\pi k_m h (p_i - p)}{q\mu B}; t_D = \frac{k_m t}{\varphi_m \mu C_{tm} r_c^2};$$

$$r_D = \frac{r}{r_c}; r_{eD} = \frac{r_e}{r_c}; \eta = \frac{k_1}{k_m}; \omega = \frac{k_m h^2 \varphi_1 C_{t1}}{k_1 r_c^2 \varphi_m C_{tm}};$$

$$C_D = \frac{C_L}{2\varphi_m C_{tm}}; z_D = \frac{z}{h}; h_{1D} = \frac{h_1}{h}; h_{2D} = \frac{h_2}{h};$$

$$\delta = \frac{r_c^2 k_1}{2k_m h_1^2}; g_D = \frac{\rho g \pi r_c^2 k_1}{qB\mu};$$

其中, p_D 为无因次压力; t_D 为无因次时间; r_D 为无因次半径; η 为等效均匀介质渗透率与基质渗透率比值; ω 反映等效均匀介质与基质系统能量比值大小; C_D 反映溶洞未充填部分与基质系统能量比值大小; z_D 为无因次厚度; δ 反映等效均匀介质与溶洞未充填区域流动比值大小; g_D 为无因次重力, 反映了重力影响.

将式(1)进行无因次变换, 可得:

$$\begin{cases} \frac{\partial^2 p_{mD}}{\partial r_D^2} + \frac{1}{r_D} \frac{\partial p_{mD}}{\partial r_D} = \frac{\partial p_{mD}}{\partial t_D}; \\ \frac{\partial^2 p_{1D}}{\partial z_D^2} = \omega \frac{\partial p_{1D}}{\partial t_D} \Big|_{r_{eD} \leq r_D \leq r_{eD} + h_{2D} \leq z_D \leq 0}; \\ p_{mD} \Big|_{r_D=0} = p_{1D} \Big|_{r_D=0} = p_{vD} \Big|_{r_D=0} = 0; \frac{\partial p_{mD}}{\partial r_D} \Big|_{r_D=r_{eD}} = 0; \\ \frac{\partial p_{1D}}{\partial z_D} \Big|_{z_D=h_{1D}} = 0; p_{mD} \Big|_{r_D=1} = p_{vD} \Big|_{r_D=1} = p_{1D} \Big|_{z_D=0} = \\ p_{vD} \Big|_{z_D=0}; \\ \frac{\partial p_{mD}}{\partial r_D} \Big|_{r_D=1} - C_D \frac{\partial p_{vD}}{\partial t_D} \Big|_{r_D=1} + \delta \frac{\partial p_{1D}}{\partial z_D} \Big|_{z_D=0} - C_D \frac{\partial p_{vD}}{\partial t_D} \Big|_{z_D=0} = \\ \frac{-1-g_D}{h_{1D}}. \end{cases} \quad (2)$$

将方程组(2)进行 Laplace 变换, 可得:

$$\begin{cases} \frac{d^2 \bar{p}_{mD}}{dr_D^2} + \frac{1}{r_D} \frac{d\bar{p}_{mD}}{dr_D} = s \bar{p}_{mD}; \\ \frac{d^2 \bar{p}_{1D}}{dz_D^2} = \omega s \bar{p}_{1D}; \\ \frac{d\bar{p}_{mD}}{dr_D} \Big|_{r_D=r_{eD}} = 0; \frac{d\bar{p}_{1D}}{dz_D} \Big|_{z_D=h_{1D}} = 0; \\ \bar{p}_{mD} \Big|_{r_D=1} = \bar{p}_{vD} \Big|_{r_D=1}; \bar{p}_{1D} \Big|_{z_D=0} = \bar{p}_{vD} \Big|_{z_D=0}; \\ \frac{d\bar{p}_{mD}}{dr_D} \Big|_{r_D=1} - C_D s \bar{p}_{vD} \Big|_{r_D=1} + \delta \frac{d\bar{p}_{1D}}{dz_D} \Big|_{z_D=0} - \\ C_D s \bar{p}_{vD} \Big|_{z_D=0} = \frac{-1-g_D}{sh_{1D}}. \end{cases} \quad (3)$$

通过计算, 可得到公式(3)在 Laplace 空间的解为:

$$\begin{cases} \bar{p}_{mD} = AI_0(r_D \sqrt{s}) + BK_0(r_D \sqrt{s}); \\ \bar{p}_{1D} = Ce^{z_D \sqrt{\omega s}} + De^{-z_D \sqrt{\omega s}}; \end{cases} \quad (4)$$

其中:

$$\begin{cases} B = \frac{-1-g_D}{sh_{1D} [\sqrt{s}M - C_D sN + E \cdot F]}; \\ A = B [K_1(r_{eD} \sqrt{s}) / I_1(r_{eD} \sqrt{s})]; \\ C = B \cdot E; D = C \cdot \exp(2h_{1D} \sqrt{\omega s}); \\ M = \frac{K_1(r_{eD} \sqrt{s})}{I_1(r_{eD} \sqrt{s})} I_0(\sqrt{s}) - K_1(\sqrt{s}); \\ N = \frac{K_1(r_{eD} \sqrt{s})}{I_1(r_{eD} \sqrt{s})} I_0(\sqrt{s}) + K_0(\sqrt{s}); \\ E = \frac{N}{1 + e^{2h_{1D} \sqrt{\omega s}}}; \\ F = \delta \sqrt{\omega s} [1 - e^{2h_{1D} \sqrt{\omega s}}] - C_D s [1 + e^{2h_{1D} \sqrt{\omega s}}]. \end{cases}$$

根据公式(4)可知, 溶洞未充填区域无因次井底

压力在 Laplace 空间解为：

$$\bar{p}_{wD} = AI_0(\sqrt{s}) + BK_0(\sqrt{s}) . \quad (5)$$

基岩—溶洞未充填区域和充填物—溶洞未充填区域窜流量分别为：

$$\bar{q}_{cD1} = -h_{1D} \frac{\partial \bar{p}_{mD}}{\partial r_D} \Big|_{r_D=1} = -h_{1D} \sqrt{s} [AI_1(\sqrt{s}) - BK_1(\sqrt{s})] , \quad (6)$$

$$\bar{q}_{cD2} = -\delta h_{1D}^2 \frac{d\bar{p}_{1D}}{dz_D} \Big|_{z_D=0} - \frac{g_D}{s} = -\delta h_{1D}^2 \sqrt{\omega s} (C - D) - \frac{g_D}{s} . \quad (7)$$

根据式(6)和式(7)，可以得到大尺度溶洞无因次流量为：

$$\bar{q}_{cD} = \bar{q}_{cD1} + \bar{q}_{cD2} . \quad (8)$$

根据 Stehfest 数值反演方法可将 Laplace 空间解转换到实空间解，最终可得到基岩—溶洞未充填区域窜流量、充填物—溶洞未充填区域窜流量和大尺度溶洞无因次流量。

2 井钻遇大尺度部分充填溶洞流动特征

2.1 流动过程分析

图 2 选取参数为： $h_{2D} = 0.5$ 、 $\omega = 5$ 、 $C_D = 20$ 、 $r_{cD} = 60$ 、 $\eta = 15.8$ 、 $\delta = 100$ 、 $g_D = 1 \times 10^{-4}$ ，通过计算可以得到井钻遇大尺度部分充填溶洞流动特征曲线，结果表明流动过程可以分为如下 4 个阶段：(1) 窜流前期(基质中流体向溶洞未充填部分窜流量 q_{cD1} 较大，溶洞充填部分流体向溶洞未充填部分窜流量 q_{cD2} 较小，窜流量 q_{cD1} 和 q_{cD2} 均随时间增加而线

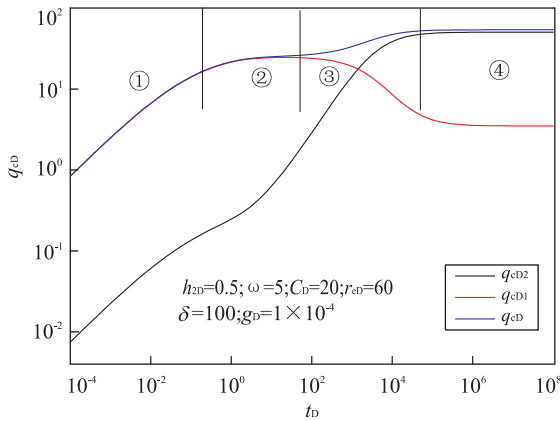


图 2 流动特征曲线

Fig.2 Typical crossflow characteristic curve for the model

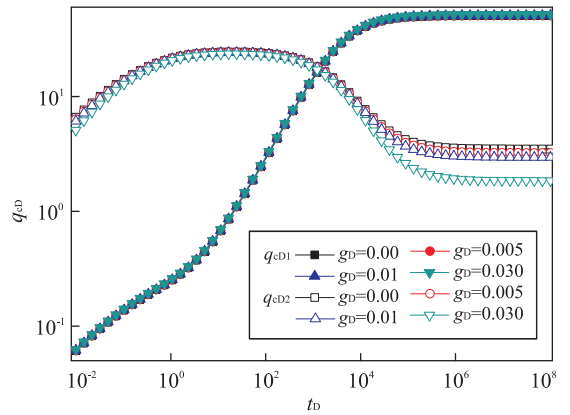


图 3 无因次重力对窜流特征曲线影响

Fig.3 Effect of dimensionless gravity on the crossflow characteristic curve

性增加，窜流特征曲线上出现第一个上升段)；(2) 窜流中前期(随着流动不断进行，基质系统与溶洞未充填部分的压差逐步减小，基质—溶洞未充填部分流动达到拟稳态，此时基质系统中流体向溶洞未充填部分窜流速度为常数， q_{cD1} 仍然大于 q_{cD2} ，窜流特征曲线上出现第一个平缓段)；(3) 窜流中后期(基质系统中流体向溶洞未充填部分窜流速度进一步降低，溶洞充填部分流体向溶洞未充填部分窜流量 q_{cD2} 超过 q_{cD1} ，窜流特征曲线上出现第二个上升段)；(4) 窜流后期(整个系统流动达到拟稳态，窜流特征曲线上出现第二个平缓段)。

2.2 敏感性分析

2.2.1 重力对窜流特征曲线影响 大尺度溶洞(洞穴)中，重力是影响充填物中流体向溶洞未充填部分窜流的重要参数(李鹤和李允, 2010; Cai *et al.*, 2012)。图 3 中选取参数分别为： $h_{2D} = 0.5$ 、 $\omega = 5$ 、 $C_D = 20$ 、 $r_{cD} = 60$ 、 $\eta = 15.8$ 、 $\delta = 100$ ，结果表明，基质—溶洞未充填部分窜流不受到重力影响，而溶洞充填物向溶洞未充填部分窜流量随着无因次重力 g_D 增大而减小，说明重力抑制了溶洞充填物中流体向溶洞未充填部分流动。主要原因是在重力影响下，溶洞充填物中流体向溶洞未充填部分窜流速度降低，造成窜流量减小。

2.2.2 无因次参数 C_D 对窜流特征曲线影响 无因次参数 C_D 反映了溶洞未充填部分与基质系统能量比值，而根据式(6)和式(7)可知， C_D 是影响充填物(或基质系统)中流体向溶洞未充填部分流动特征的重要因素(图 4)。图 4 中选取参数分别为： $h_{2D} = 0.5$ 、 $\omega = 5$ 、 $r_{cD} = 60$ 、 $\delta = 100$ 、 $\eta = 15.8$ 、 $g_D = 1 \times 10^{-4}$ ，结果表明无因次参数 C_D 主要影响窜流初期，基岩—溶

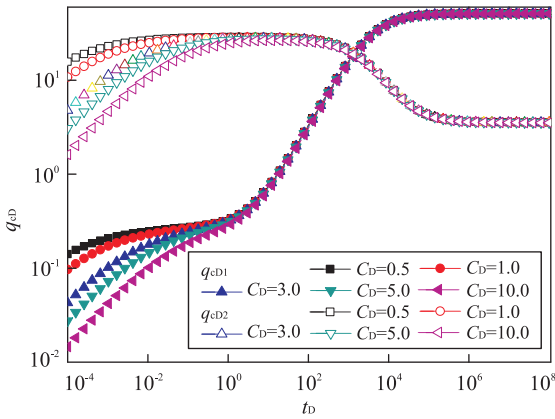


图 4 无因次变量 C_D 对窜流特征曲线影响

Fig.4 Effect of dimensionless variable C_D on the crossflow characteristic curve

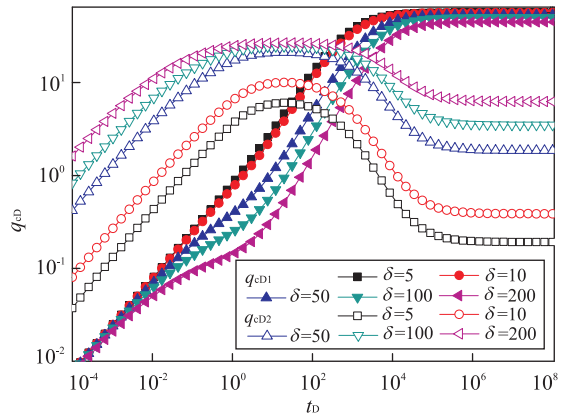


图 6 无因次变量 δ 对窜流特征曲线影响

Fig.6 Effect of dimensionless variable δ on the crossflow characteristic curve

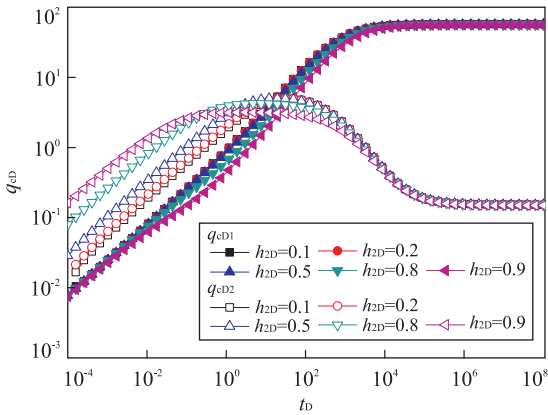


图 5 无因次变量 h_{2D} 对窜流特征曲线影响

Fig.5 Effect of dimensionless variable h_{2D} on the crossflow characteristic curve

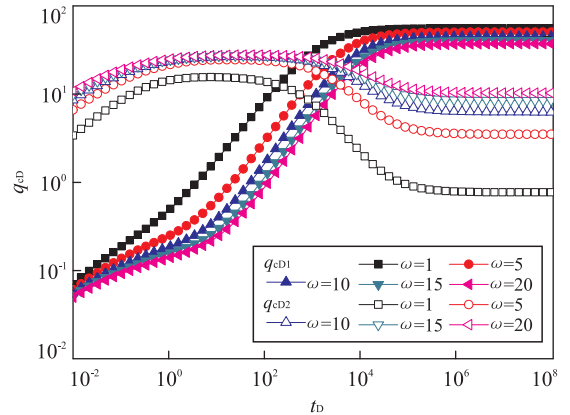


图 7 无因次变量 ω 对窜流特征曲线影响

Fig.7 Effect of dimensionless variable ω on the crossflow characteristic curve

洞未充填区域和充填物—溶洞未充填区域窜流量均随着无因次参数 C_D 增大而减小,主要原因是无因次参数 C_D 越大,溶洞未充填部分能量越大,溶基质系统和溶洞未充填部分能量差越小,基质系统中流体向溶洞未充填部分窜流量越小。

2.2.3 无因次参数 h_{2D} 对窜流特征曲线影响 无因次参数 h_{2D} 反映了溶洞充填程度大小,且溶洞充填程度对窜流特征曲线影响如图 5 所示,其中选取参数分别为: $\omega=5, C_D=20, r_{cD}=60, \delta=100, \eta=15.8, g_D=1 \times 10^{-4}$, 结果表明无因次参数 h_{2D} 越大,充填物—溶洞未充填区域窜流量 q_{cD2} 越大,而基质中流体向溶洞未充填部分窜流量 q_{cD1} 越小。主要原因是无因次参数 h_{2D} 越大,无因次参数 h_{1D} 越小,溶洞充填程度越大,充填物能量越大,充填物中流体向溶洞未充填区域窜流量越大,基质中流体向溶洞未充填部分窜流量越小。

2.2.4 无因次参数 δ 和 ω 对窜流特征曲线影响

图 6 为无因次参数 δ 对窜流特征曲线影响,其中选取参数分别为: $h_{2D}=0.5, \omega=5, C_D=20, r_{cD}=60, \eta=15.8, g_D=1 \times 10^{-4}$, 结果表明基岩—溶洞未充填区域窜流量 q_{cD1} 随着无因次参数 δ 增大而减小,而溶洞充填物—溶洞未充填区域窜流量 q_{cD2} 随着无因次参数 δ 增大而增大。主要原因是无因次参数 δ 越大,溶洞充填物—溶洞未充填区域流动能力越强,而基质—溶洞未充填区域流动能力越弱。

图 7 为无因次参数 ω 对窜流特征曲线影响,其中选取参数分别为: $h_{2D}=0.5, C_D=20, r_{cD}=60, \delta=100, \eta=15.8, g_D=1 \times 10^{-4}$, 结果表明基岩—溶洞未充填区域窜流量 q_{cD1} 随着无因次参数 ω 增大而减小,而溶洞充填物—溶洞未充填区域窜流量 q_{cD2} 随着无因次参数 ω 增大而增大。主要原因是无因次参数 ω 越大,溶洞充填物能量越大,而基质系统

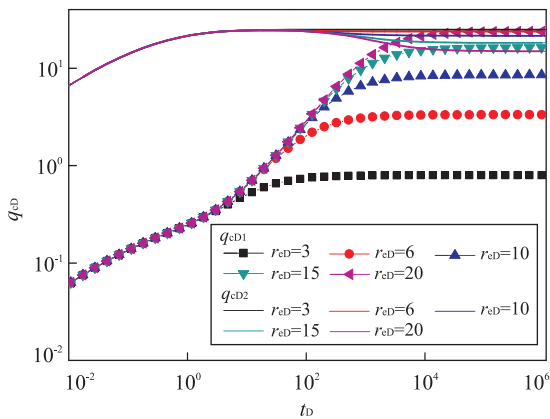


图 8 无因次变量 r_{eD} 对窜流特征曲线影响

Fig.8 Effect of dimensionless variable r_{eD} on the crossflow characteristic curve

能量越小。

2.2.5 无因次参数 r_{eD} 对窜流特征曲线影响 图 8 为无因次参数 r_{eD} 对窜流特征曲线影响,其中选取参数分别为: $h_{2D} = 0.5$ 、 $\omega = 5$ 、 $C_D = 20$ 、 $\delta = 100$ 、 $\eta = 15.8$ 、 $g_D = 1 \times 10^{-4}$,结果表明无因次参数 r_{eD} 主要影响窜流中后期,且无因次参数 r_{eD} 越大,基岩—溶洞未充填区域窜流量 q_{cD1} 越大,溶洞充填物—溶洞未充填区域窜流量 q_{cD2} 越小。主要原因是无因次参数 r_{eD} 越大,基质系统能量越大,充填物能量越小。

3 结论

(1) 基于缝洞型油藏大尺度溶洞充填特征,本研究建立了井钻遇大尺度部分充填溶洞数学模型,得到了基岩—溶洞未充填区域窜流量、充填物—溶洞未充填区域窜流量和溶洞无因次流量。

(2) 溶洞窜流特征曲线可划分为 4 个流动阶段。其中前 2 个阶段主要为基质中流体向溶洞未充填部分窜流,而后 2 个阶段主要为溶洞充填物流体向溶洞未充填部分窜流。在流动前期,基质和未充填溶洞间流体交换对部分充填溶洞流量贡献较大;而流动后期,溶洞充填物和未充填溶洞间流体交换对部分充填洞流量贡献较大。

(3) 整个窜流过程主要受到重力、溶洞未充填部分能量、溶洞充填程度和基质系统能量等因素影响。基质—溶洞未充填部分窜流不受到重力影响,重力抑制充填物中流体向溶洞未充填部分流动;溶洞充填物(或基质系统)和溶洞未充填部分能量差越小,溶洞充填物(或基质系统)中流体向溶洞未充填部分窜流量越小。

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